

Plan of Lectures

- I. Standard Neutrino Properties and Mass Terms (Beyond Standard)**
- II. Effects of ν Mass: Neutrino Oscillations (Vacuum)**
- III. Matter Effects in Neutrino Oscillations**
- IV. The Emerging Picture and Some Lessons**

Summary I+II+III

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- **Atmospheric, K2K and MINOS** (+ negative SBL searches)
 - $\Rightarrow \nu_\mu \rightarrow \nu_\tau$ with $\Delta m^2 \sim 2 \times 10^{-3} \text{ eV}^2$ and $\tan^2 \theta \sim 1$
- **Solar and KamLAND**
 - $\Rightarrow \nu_e \rightarrow \nu_\mu, \nu_\tau$ with $\Delta m^2 \sim 8 \times 10^{-5} \text{ eV}^2$ and $\tan^2 \theta \sim 0.4$

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- Can we fit all together? What can we learn from all this?

Answer: Today

Plan of Lecture IV

Emerging Picture and Some Lessons

3ν Oscillations

Some Lessons:

The Need of New Physics

The Possibility of Leptogenesis

- We have learned:

- * Atmospheric ν_μ disappear ($> 15\sigma$) most likely to ν_τ
- * K2K: accelerator ν_μ disappear at $L \sim 250$ Km with E -distortion ($\sim 2.5\text{--}4\sigma$)
- * MINOS: accelerator ν_μ disappear at $L \sim 735$ Km with E -distortion ($\sim 5\sigma$)
- * Solar ν_e convert to ν_μ or ν_τ ($> 7\sigma$)
- * KamLAND: reactor $\bar{\nu}_e$ disappear at $L \sim 200$ Km with E -distortion ($\gtrsim 3\sigma$ CL)
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- We have important information (mostly constraints) from:

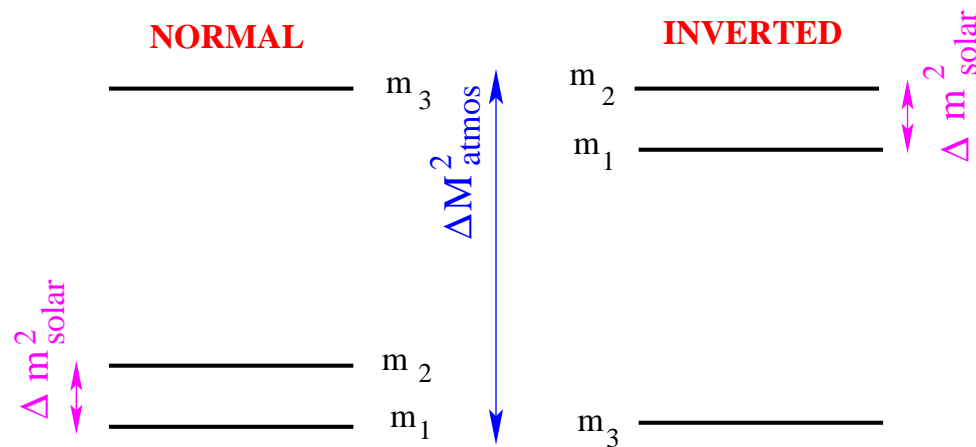
- * The line shape of the Z: $N_{\text{weak}} = 3$
- * Limits from Short Distance Oscillation Searches at Reactor and Accelerators
- * Direct mass measurements: ${}^3\text{H} \rightarrow {}^3\text{He} + e^- + \bar{\nu}_e$ and ν -less $\beta\beta$ decay
- * From Astrophysics and Cosmology: BBN, CMBR, LSS ...

Solar+Atmospheric+Reactor+LBL 3ν Oscillations

U : 3 angles, 1 CP-phase
 + (2 Majorana phases)

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{-i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{21} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Two mass schemes



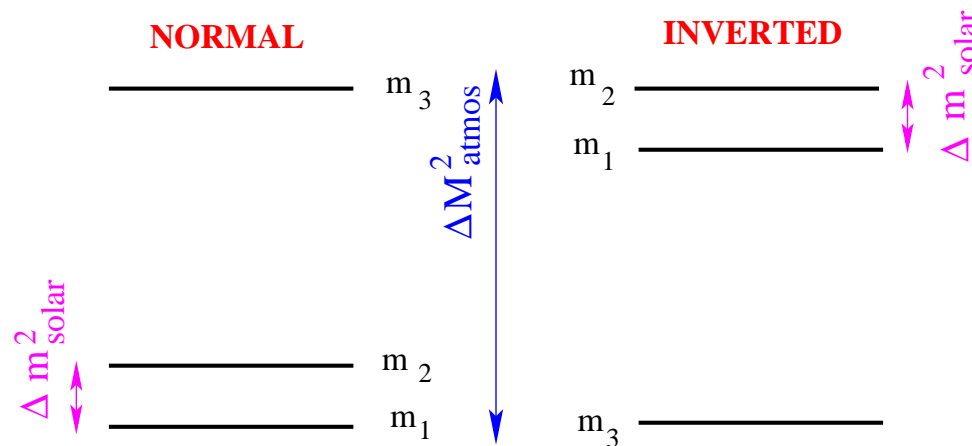
2ν oscillation analysis $\Rightarrow \Delta m_{21}^2 = \Delta m_{\odot}^2 \ll \Delta M_{\text{atm}}^2 \simeq \pm \Delta m_{32}^2 \simeq \pm \Delta m_{31}^2$

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Generic 3ν mixing effects:

- Effects due to θ_{13}
- Difference between **Inverted** and **Normal**
- Interference of **two wavelength** oscillations
- **CP violation** due to phase δ

3- ν Neutrino Oscillations

- In general one has to solve: $i \frac{d\vec{\nu}}{dt} = H \vec{\nu}$

$$H = U \cdot H_0^d \cdot U^\dagger + V$$

$$H_0^d = \frac{1}{2E_\nu} \text{diag} \left(-\Delta m_{21}^2, 0, \Delta m_{32}^2 \right)$$

$$V = \text{diag} \left(\pm \sqrt{2} G_F N_e, 0, 0 \right)$$

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- Hierarchical approximation: $\Delta m_{21}^2 \ll \Delta m_{31}^2 \sim \Delta m_{32}^2$

- * For $\theta_{13} = 0$ solar and atmospheric oscillations decouple \Rightarrow Normal \equiv Inverted
 - Solar and KamLAND $\rightarrow \Delta m_{21}^2 = \Delta m_\odot^2 \quad \theta_{12} = \theta_\odot$
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- K2K+MINOS Probabilities Independent of $\theta_{12}, \Delta m_{21}^2$ 4

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• Hierarchical approximation: $\Delta m_{21}^2 \ll \Delta m_{31}^2 \sim \Delta m_{32}^2 \Rightarrow$ neglect Δm_{21}^2 in ATM

$$P_{ee} = 1 - 4s_{13,m}^2 c_{13,m}^2 S_{31}$$

$$P_{\mu\mu} = 1 - 4s_{13,m}^2 c_{13,m}^2 s_{23}^4 S_{31} - 4s_{13,m}^2 s_{23}^2 c_{23}^2 S_{21} - 4c_{13,m}^2 s_{23}^2 c_{23}^2 S_{32}$$

$$P_{e\mu} = 4s_{13,m}^2 c_{13,m}^2 s_{23}^2 S_{31}$$

$$S_{ij} = \sin^2 \left(\frac{\Delta\mu_{ij}^2 L}{4E_\nu} \right)$$

Pantaleone 94; Used by many ...

$$\Delta\mu_{21}^2 = \frac{\Delta m_{32}^2}{2} \left(\frac{\sin 2\theta_{13}}{\sin 2\theta_{13,m}} - 1 \right) - E_\nu V_e$$

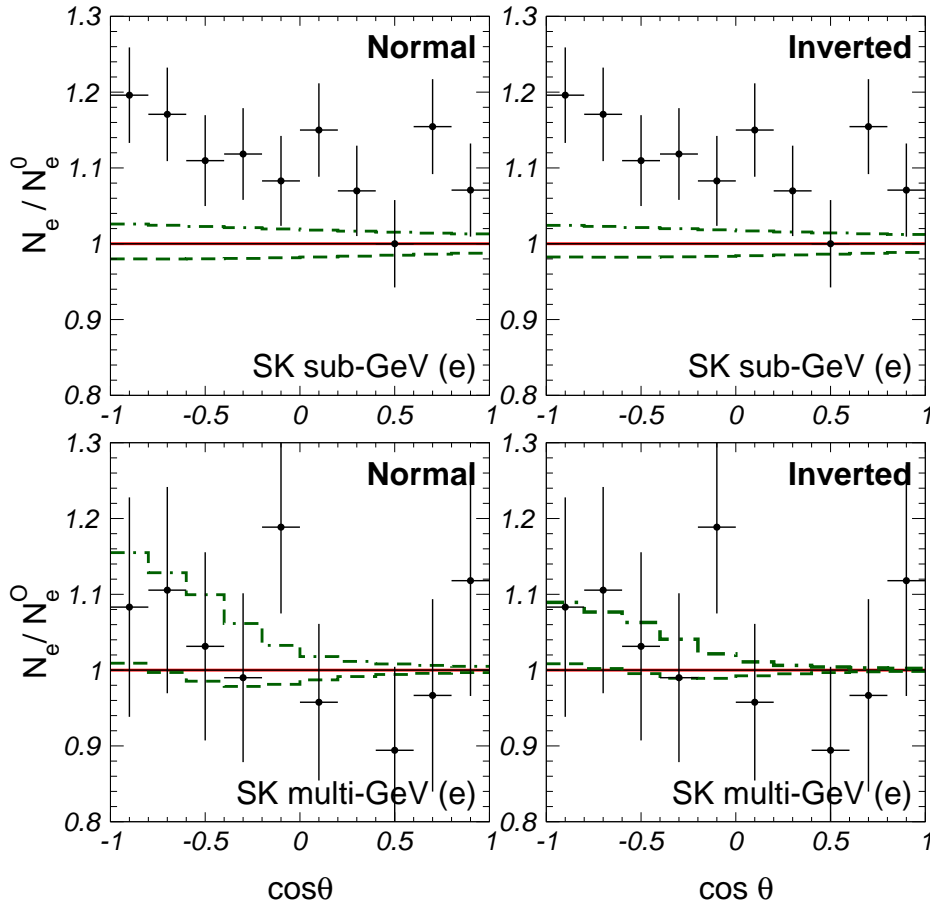
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Ahkmedov, Dighe, Lipari, Smirnov 99; Petcov, Maris 98; Palomares, Petcov, 03



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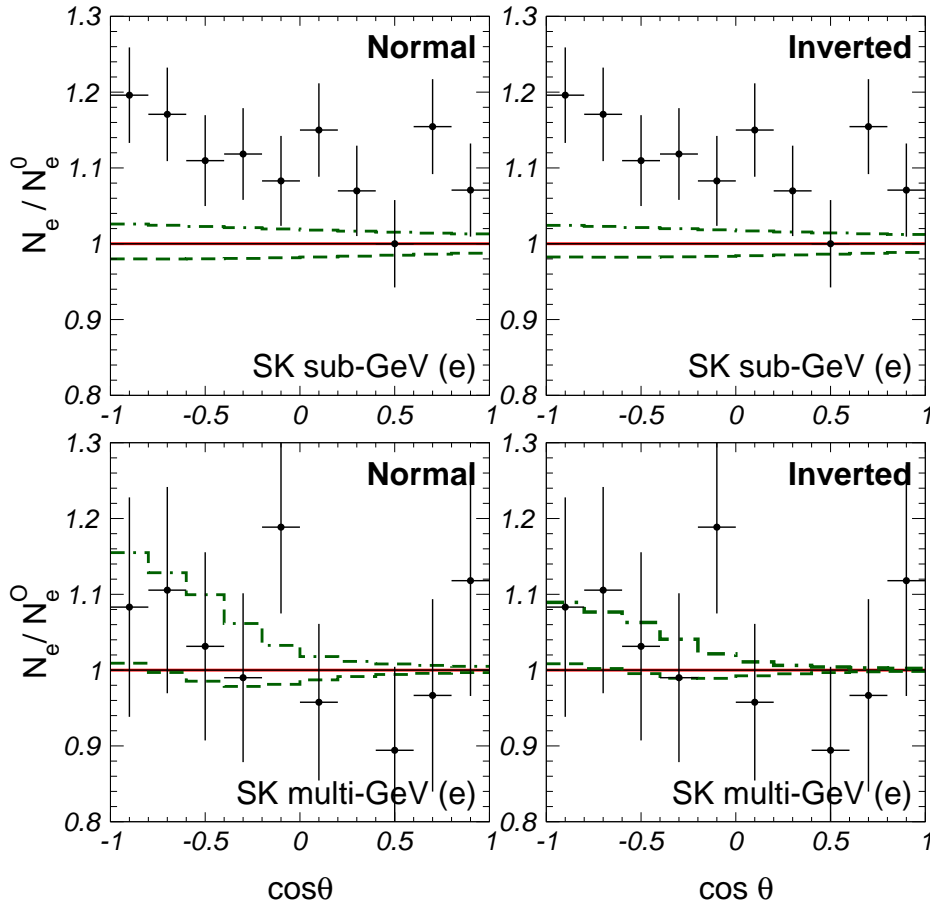
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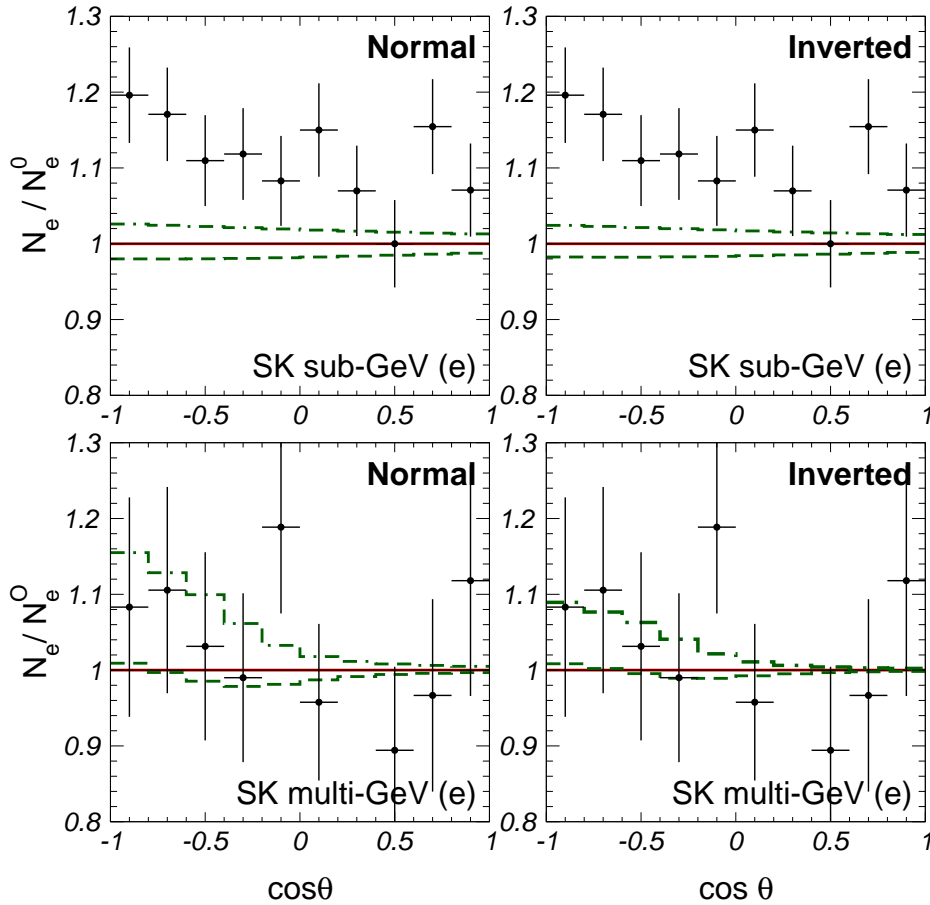
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Larger Effect in Normal

Possible Sensitivity to Mass Ordering

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Possible Sensitivity to Mass Ordering

• Sub-GeV: Vacuum Osc: Smaller Effect

$$r \simeq 2 \Rightarrow \theta_{23} < \frac{\pi}{4} \Rightarrow s_{23}^2 < \frac{1}{2} \Rightarrow N_e(\theta_{13}) < N_{e0}$$

$$\theta_{23} > \frac{\pi}{4} \Rightarrow s_{23}^2 > \frac{1}{2} \Rightarrow N_e(\theta_{13}) > N_{e0}$$

Δm_{21}^2 effects in ATM Data

Smirnov, Peres 99,01; Fogli, Lisi, Marrone 01; MC G-G, Maltoni 02; MCG-G, Maltoni, Smirnov hep-ph/0408170

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- Neglecting θ_{13} :

$$P_{ee} = 1 - P_{e2}$$

$$P_{e\mu} = c_{23}^2 P_{e2}$$

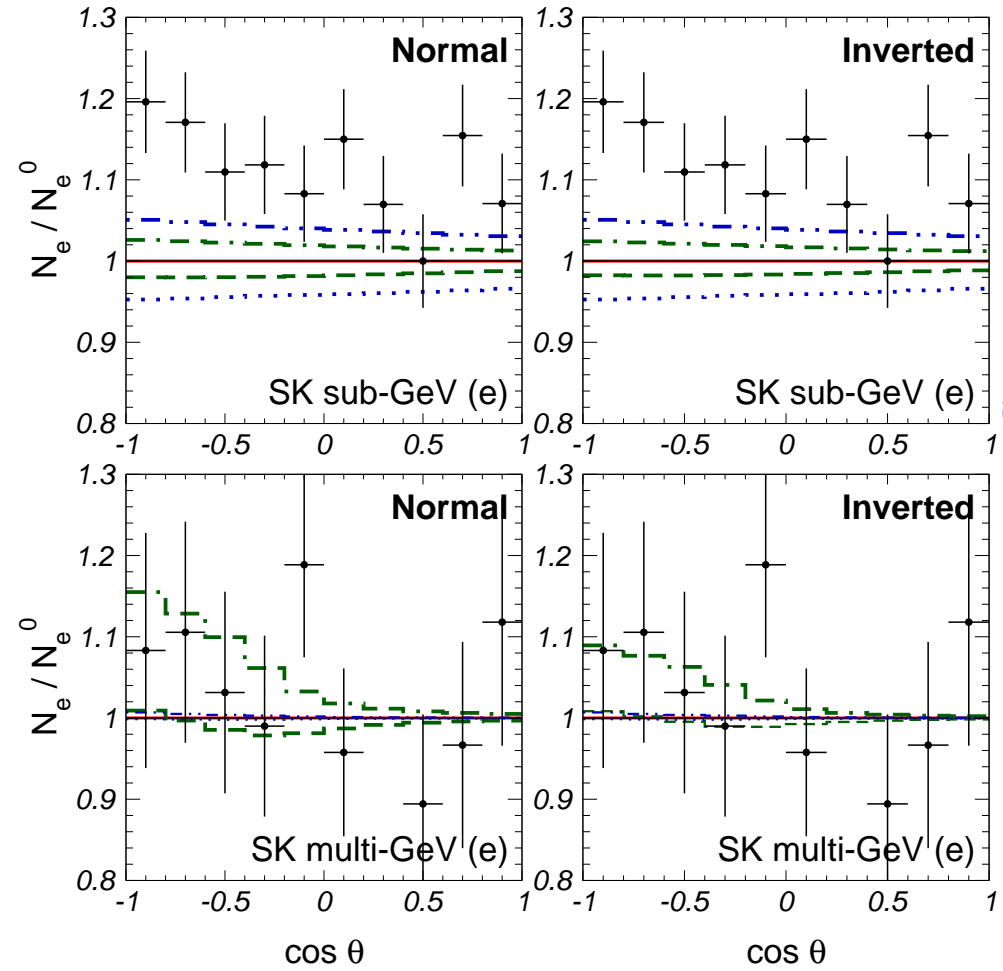
$$P_{\mu\mu} = 1 - c_{23}^4 P_{e2} - 2s_{23}^2 c_{23}^2 \left[1 - \sqrt{1 - P_{e2}} \cos \phi \right]$$

$$P_{e2} = \sin^2 2\theta_{12,m} \sin^2 \left(\frac{\Delta m_{21}^2 L}{4E_\nu} \frac{\sin 2\theta_{12}}{\sin 2\theta_{12,m}} \right)$$

$$\sin 2\theta_{12,m} = \frac{\sin 2\theta_{12}}{\sqrt{\left(\cos 2\theta_{12} \mp \frac{2E_\nu V_e}{\Delta m_{21}^2} \right)^2 + \sin^2 2\theta_{12}}}$$

$$\phi \approx \left(\Delta m_{31}^2 + s_{12}^2 \Delta m_{21}^2 \right) \frac{L}{2E_\nu}$$

Δm_{21}^2 effects in ATM Data



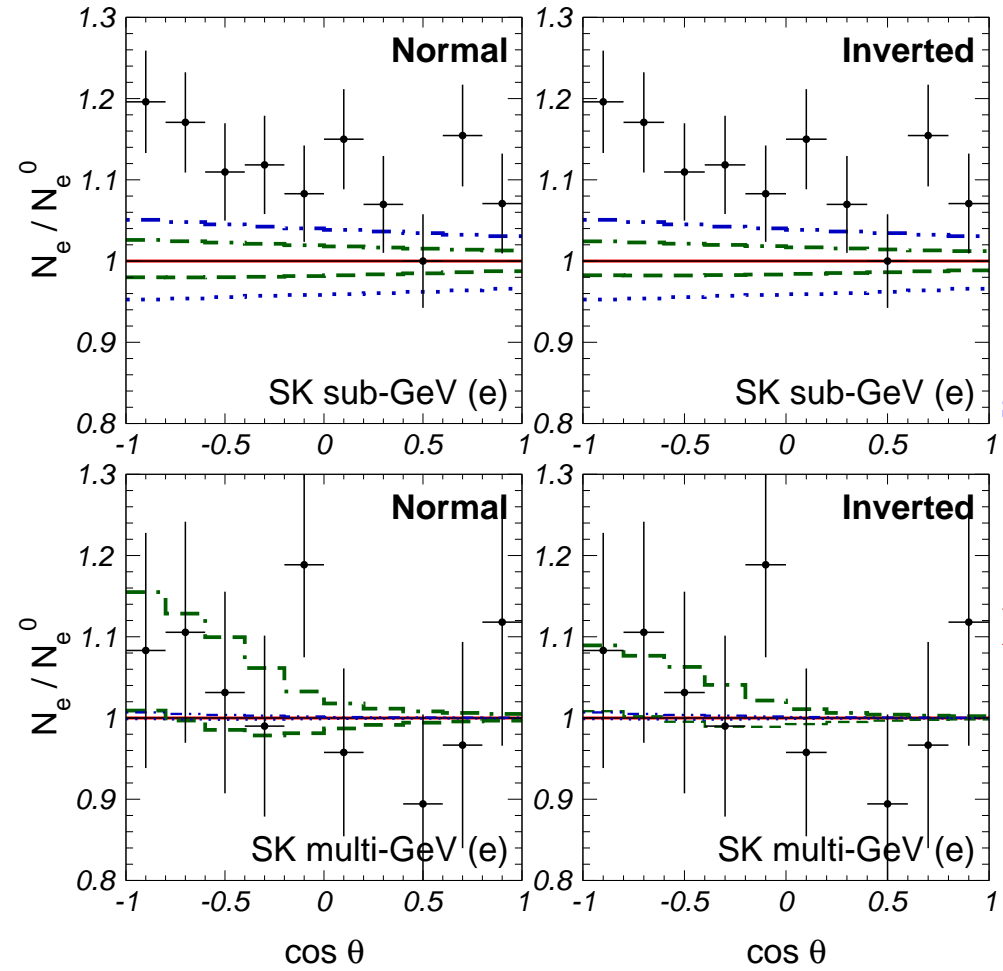
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For Sub-GeV:

$$P_{e2} = \frac{(\Delta m_{21}^2)^2}{(2EV_e)^2} \sin^2 2\theta_{12} \sin^2 \frac{V_e L}{2}$$

$$\theta_{23} < \frac{\pi}{4} \Rightarrow c_{23}^2 > \frac{1}{2} \Rightarrow N_e(\theta_{13}) > N_{e0}$$

$$\theta_{23} > \frac{\pi}{4} \Rightarrow c_{23}^2 < \frac{1}{2} \Rightarrow N_e(\theta_{13}) < N_{e0}$$

- \Rightarrow Sensitiv to Deviations from Maximal θ_{23}
- \Rightarrow Sensitivity to Octant of θ_{23}
(even for vanishing θ_{13})
- \Rightarrow Effect proportional to $(\Delta m_{21}^2)^2$

- $s_{13}^2=0.04, s_{23}^2=0.35, \Delta m_{21}^2=0$
- .- $s_{13}^2=0.04, s_{23}^2=0.65, \Delta m_{21}^2=0$
- $s_{13}^2=0.00, s_{23}^2=0.35, \Delta m_{21}^2=10^{-4} \text{ eV}^2$
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Beyond Hierarchical: Effect $\theta_{13} \times \Delta m_{21}^2$ in ATM

Smirnov, Peres 01,03, MC G-G, Maltoni 02

For sub-GeV energies

$$\frac{N_e}{N_e^0} - 1 \simeq \overline{P_{e2}} \overline{r} (c_{23}^2 - \frac{1}{\overline{r}}) + 2 \tilde{s}_{13}^2 \overline{r} (s_{23}^2 - \frac{1}{\overline{r}}) - \overline{r} \tilde{s}_{13} \tilde{c}_{13}^2 \sin 2\theta_{23} (\cos \delta_{CP} \overline{R_2} - \sin \delta_{CP} \overline{I_2})$$

$$P_{e2} = \sin^2 2\theta_{12,m} \sin^2 \frac{\phi_m}{2}$$

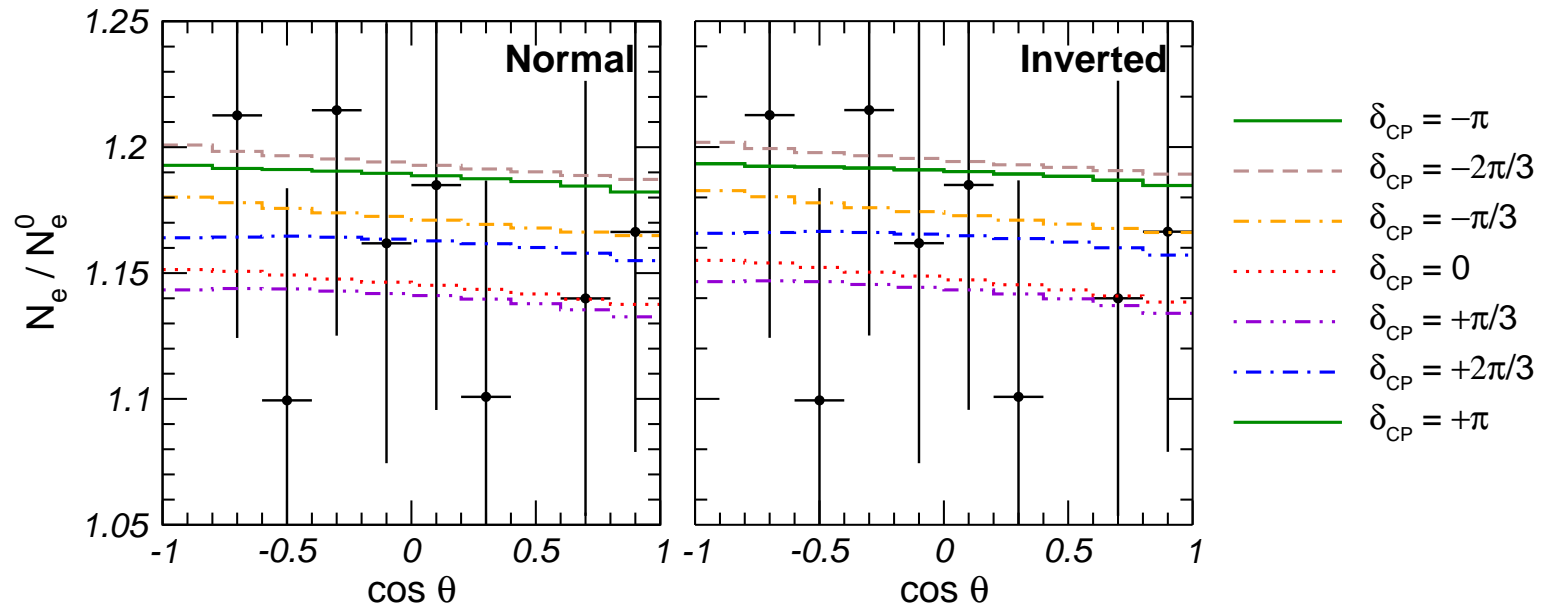
$$\sin 2\theta_{12,m} = \frac{\sin 2\theta_{12}}{\sqrt{(\cos 2\theta_{12} \mp \frac{2E_\nu V_e}{\Delta m_{21}^2})^2 + \sin^2 2\theta_{12}}}$$

$$R_2 = -\sin 2\theta_{12,m} \cos 2\theta_{12,m} \sin^2 \frac{\phi_m}{2}$$

$$I_2 = -\frac{1}{2} \sin 2\theta_{12,m} \sin \phi_m$$

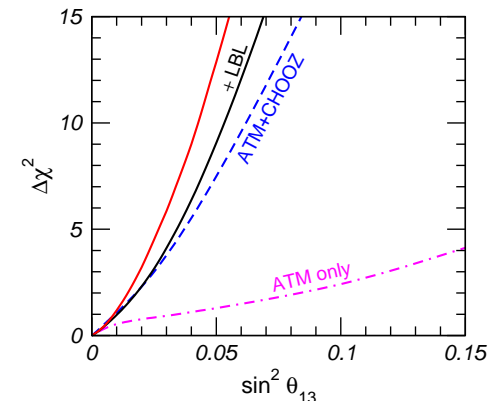
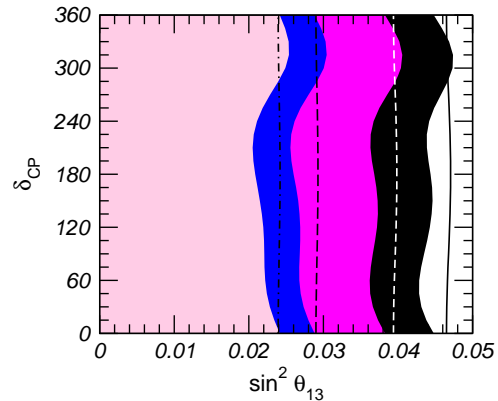
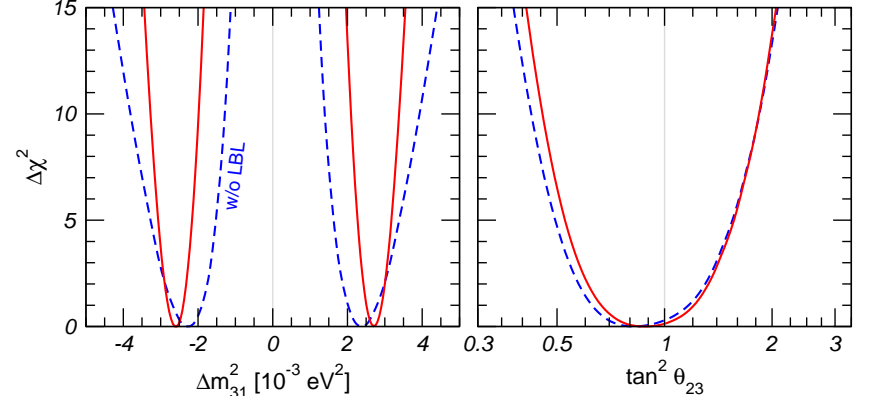
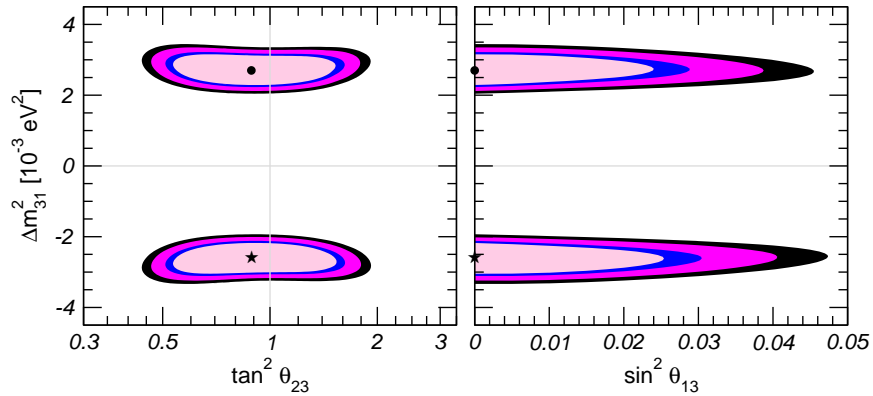
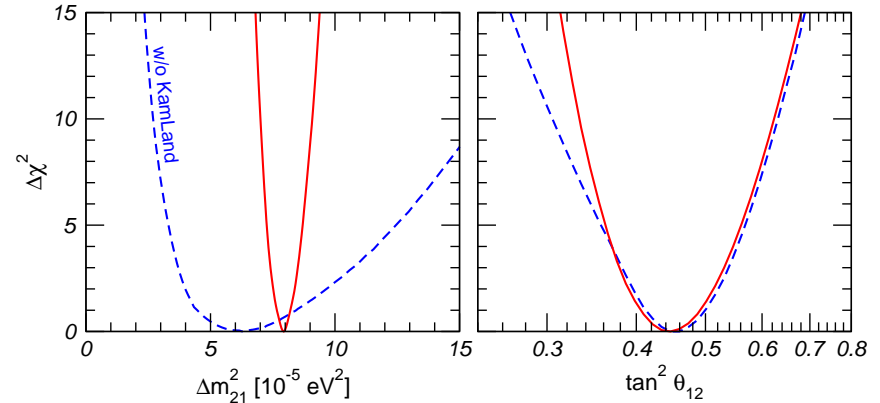
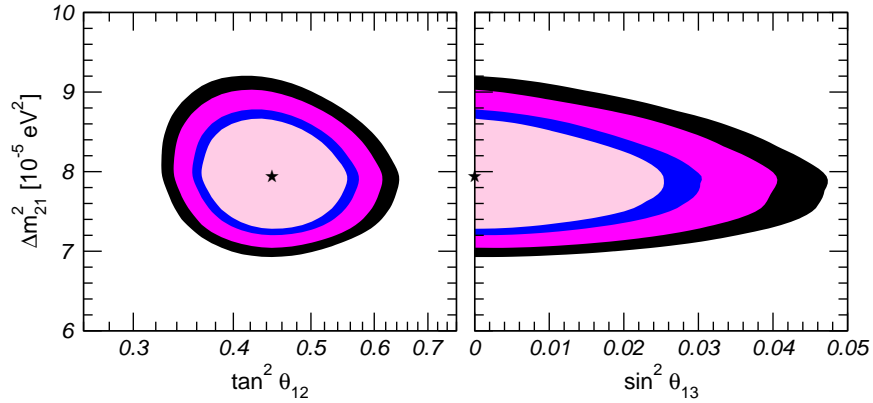
$$\tilde{\theta}_{13} \approx \theta_{13} \left(1 + \frac{2E_\nu V_e}{\Delta m_{31}^2} \right)$$

$$\phi \approx (\Delta m_{31}^2 + s_{12}^2 \Delta m_{21}^2) \frac{L}{2E_\nu}$$



Global Analysis: Three Neutrino Oscillations

M.C. G-G, M.Maltoni, ArXiV/0704.1800



Global Analysis: Three Neutrino Oscillations

The derived ranges:

$$\Delta m_{21}^2 = 7.7^{+0.22}_{-0.21} \left(\begin{array}{c} +0.67 \\ -0.61 \end{array} \right) \times 10^{-5} \text{ eV}^2 \quad \left| \Delta m_{31}^2 \right| = 2.37 \pm 0.17 (0.46) \times 10^{-3} \text{ eV}^2$$

$$|U_{LEP}|_{3\sigma} = \begin{pmatrix} 0.79 \rightarrow 0.86 & 0.50 \rightarrow 0.61 & 0.00 \rightarrow 0.20 \\ 0.25 \rightarrow 0.53 & 0.47 \rightarrow 0.73 & 0.56 \rightarrow 0.79 \\ 0.21 \rightarrow 0.51 & 0.42 \rightarrow 0.69 & 0.61 \rightarrow 0.83 \end{pmatrix}$$

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with structure

$$|U_{LEP}| \simeq \begin{pmatrix} \frac{1}{\sqrt{2}}(1 + \mathcal{O}(\lambda)) & \frac{1}{\sqrt{2}}(1 - \mathcal{O}(\lambda)) & \epsilon \\ -\frac{1}{2}(1 - \mathcal{O}(\lambda) + \epsilon) & \frac{1}{2}(1 + \mathcal{O}(\lambda) - \epsilon) & \frac{1}{\sqrt{2}} \\ \frac{1}{2}(1 - \mathcal{O}(\lambda) - \epsilon) & -\frac{1}{2}(1 + \mathcal{O}(\lambda) - \epsilon) & \frac{1}{\sqrt{2}} \end{pmatrix} \quad \begin{matrix} \lambda \sim 0.2 \\ \epsilon \lesssim 0.2 \end{matrix}$$

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very different from quark's

$$|U_{CKM}| \simeq \begin{pmatrix} 1 & \mathcal{O}(\lambda) & \mathcal{O}(\lambda^3) \\ \mathcal{O}(\lambda) & 1 & \mathcal{O}(\lambda^2) \\ \mathcal{O}(\lambda^3) & \mathcal{O}(\lambda^2) & 1 \end{pmatrix} \quad \lambda \sim 0.2$$

Open Questions

We still ignore:

- (1) Is $\theta_{13} \neq 0$? How small?
- (2) Is $\theta_{23} = \frac{\pi}{4}$? If not, is it $>$ or $<$?
- (3) Is there CP violation in the leptons (is $\delta \neq 0, \pi$)?
- (4) What is the ordering of the neutrino states?
- (5) Are neutrino masses:
 - hierarchical: $m_i - m_j \sim m_i + m_j$?
 - degenerated: $m_i - m_j \ll m_i + m_j$?
- (6) Dirac or Majorana?

To answer (1)–(4): Proposed new generation ν osc experiments:

- Medium Baseline Reactor Experiment: *Double-Chooz, Daya Bay*
- Conventional (=from π decay) Superbeams: *T2K, Nova (?)*
- ν -factory: clean ν beam from μ decay
- ν_e or $\bar{\nu}_e$ beam from nuclear β decay (β beam)

Some Lessons: New Physics

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Thus the most striking implication of ν masses:

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And it is also the only solid evidence!

To go further one has to be cautious...

Lessons: The Scale of New Physics

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If SM is an effective low energy theory, for $E \ll \Lambda_{\text{NP}}$

- The same particle content as the SM and same pattern of symmetry breaking
- But there can be non-renormalizable
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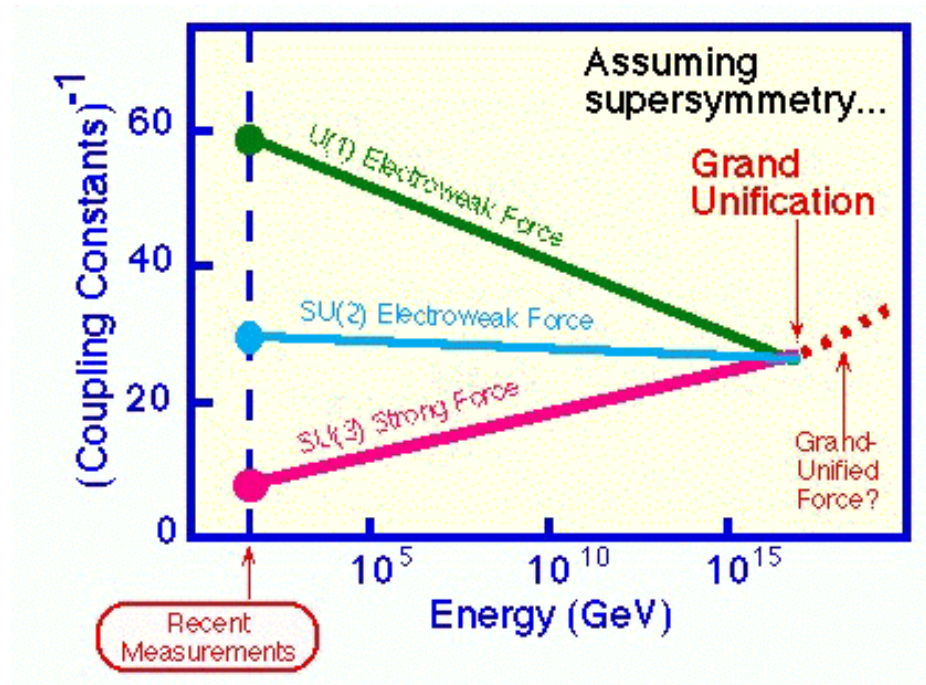
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But this is scale was already known to particle physicists...

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$$m_\nu > \sqrt{\Delta m_{\text{atm}}^2} \sim 0.05 \text{eV} \Rightarrow 10^{10} < \Lambda_{\text{NP}} < 10^{15} \text{GeV}$$

New Physics Scale close to Grand Unification scale



Also the generated neutrino mass term is Majorana :
 \Rightarrow It violates total lepton number

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Lessons:

– \mathcal{L}_{NP} contains 18 parameters which we want to know

– \mathcal{L}_5 contains 9 parameters which we can measure

\Rightarrow Same \mathcal{O}_5 can give very different \mathcal{L}_{NP}

\Rightarrow It is *difficult* to “imply” bottom-up (model independently)

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Baryogenesis and the SM

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Leptogenesis \equiv generation of lepton asymmetry Y_L

• At the electroweak transition sphaleron processes:

$$\Rightarrow Y_L \text{ is transformed in } Y_B \simeq -\frac{Y_L}{2}$$

- In the the **See-saw** mechanism

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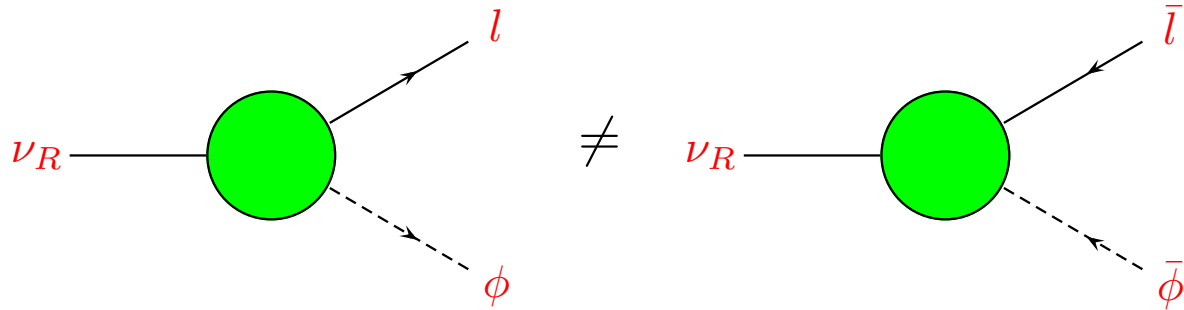
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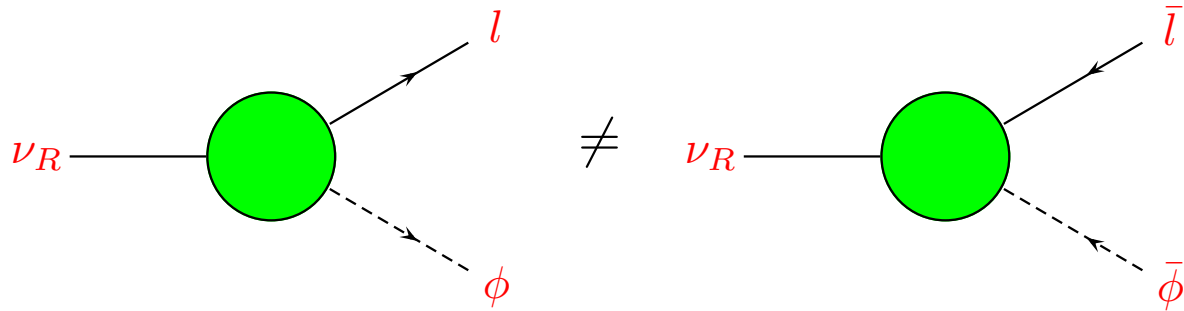
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$$\Rightarrow |\epsilon_L| \lesssim 0.1 \frac{M_{\nu_{R1}}}{\langle \phi \rangle^2} (m_{\nu_3} - m_{\nu_1})$$

$$Y_L = \frac{n_{\nu_R}}{s} \epsilon_L d \sim 10^{-3} d \epsilon_L$$

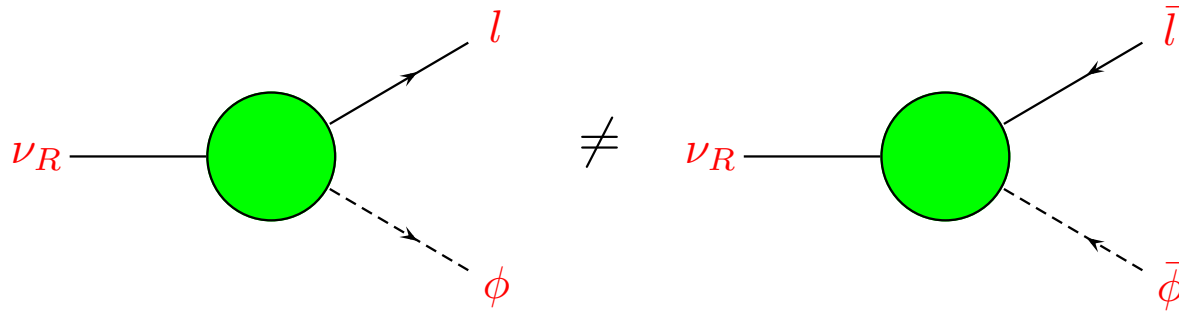
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Out of Equilibrium condition $\Gamma_{\nu_R} \ll H|_{T=M_{\nu_R}} \Rightarrow \tilde{m}_1 \equiv \frac{(\lambda \lambda^\dagger)_{11} \langle \phi \rangle^2}{M_{\nu_{R1}}} \lesssim 5 \times 10^{-3} eV$

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$$M^\nu = \begin{pmatrix} 0 & m_D \\ m_D^T & M_R \end{pmatrix}$$

$m_D = \lambda \langle \phi \rangle$ is a 3×3 matrix

M_R is a 3×3 symmetric matrix

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$\Rightarrow M^\nu$ has 6 physical phases

\Rightarrow It is easy to generate $\epsilon_L \sim 10^{-6}$

$\Rightarrow m_{\text{light}}^\nu = m_D^T M_N^{-1} m_D$ has 3 physical phases

Oscillation experiments can only see one of these three phases

\Rightarrow No direct correspondence between CPV in leptogenesis and CPV in oscillations

- The final Y_B depends on:
 - ϵ_L the CP asymmetry
 - $M_{\nu_{R1}}$ the mass of the lightest ν_R
 - $\tilde{m}_1 \equiv \frac{(\lambda\lambda^\dagger)_{11}^2 \langle \phi \rangle^2}{M_{\nu_{R1}}}$ the *effective* neutrino mass
 - $m_{\nu_1}^2 + m_{\nu_2}^2 + m_{\nu_3}^2$ the sum of the light neutrinos mass squared
- To generate the required Y_B :
 - $M_{\nu_{R1}} \gtrsim 4 \times 10^8 \text{ GeV}$
 - $m_{\nu_3} \lesssim 0.12 \text{ eV}$
 - Large CP phases
 - The CP violating phase relevant for leptogenesis may not be the same as the one relevant for oscillations

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- Neutrino oscillation searches have shown us

– $\Delta m_{31}^2 \sim 2 \times 10^{-3} \text{ eV}^2$ and $\Delta m_{21}^2 \sim 8 \times 10^{-5} \text{ eV}^2 \Rightarrow \nu$'s are massive

$$-|U_{\text{LEP}}| \simeq \begin{pmatrix} \frac{1}{\sqrt{2}}(1 + \mathcal{O}(\lambda)) & \frac{1}{\sqrt{2}}(1 - \mathcal{O}(\lambda)) & \epsilon \\ -\frac{1}{2}(1 - \mathcal{O}(\lambda) + \epsilon) & \frac{1}{2}(1 + \mathcal{O}(\lambda) - \epsilon) & \frac{1}{\sqrt{2}} \\ \frac{1}{2}(1 - \mathcal{O}(\lambda) - \epsilon) & -\frac{1}{2}(1 + \mathcal{O}(\lambda) - \epsilon) & \frac{1}{\sqrt{2}} \end{pmatrix} \Rightarrow \begin{array}{l} \lambda \sim 0.2 \\ \epsilon \lesssim 0.2 \\ \text{Different from } U_{CKM} \end{array}$$

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- breaking* total lepton number \rightarrow Majorana $\nu : \nu = \nu^C$
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(a) *breaking* total lepton number \rightarrow Majorana $\nu : \nu = \nu^C$

(b) *conserving* total lepton number \rightarrow Dirac $\nu : \nu \neq \nu^C$

- Majorana ν 's are more *Natural*: appear generically if SM is a LE effective theory

$$- \Lambda_{NP} \lesssim 10^{15} \text{ GeV}$$

– Results Fit well with GUT expectations

– Leptogenesis may explain the baryon asymmetry

Conclusions

- Still open questions

Is $\theta_{13} \neq 0$?

Is there CP violation in the leptons (is $\delta \neq 0, \pi$)?

Is θ_{23} large or maximal?

Normal or Inverted mass ordering?

Are neutrino masses:

hierarchical: $m_i - m_j \sim m_i + m_j$?

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Also no-oscillation experiments:

- ν -less $\beta\beta$ decay, ^3H beta decay
- Interesting input from cosmological data

Rich and Challenging Experimental Program